Infinitely divisible distributions and the Lévy-Khintchine formula

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Aditya Vaidyanathan Lévy-Khintchine formula

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Some definitions

Let X be a real-valued random variable with law μ_X . Recall that X is said to be infinitely divisible if for every $n \in \mathbb{N}$ we can find i.i.d. random variables $Y_1^{(n)}, \ldots, Y_n^{(n)}$ such that

$$X\stackrel{\mathsf{d}}{=}Y_1^{(n)}+\cdots+Y_n^{(n)},$$

or equivalently if there exists a probability measure μ_Y (depending on *n*) such that

$$\mu_X = \underbrace{\mu_Y * \cdots * \mu_Y}_{} := (\mu_Y)^{n*}.$$

n times

Some examples are

- the Gaussian distribution.
- the compound Poisson distribution.
- the Dirac delta distribution.
- stable distributions.

Some definitions

A real-valued stochastic process $\{X_t\}_{t\geq 0}$ is said to be a Lévy process if:

1.
$$X_0 = 0$$
 a.s.

- 2. it has stationary and independent increments.
- 3. it is stochastically continuous, i.e. for any a > 0,

$$\lim_{t\to s} P(|X_t-X_s|>a)=0.$$

Every Lévy process has a version which is càdlàg. Some examples of Lévy processes are:

- if $X_t \sim N(0, t)$ then we have a standard Brownian motion.
- if $X_t = \sum_{k=1}^{N_t} \xi_k$, where $N_t \sim \text{Poiss}(\lambda t)$ and $\{\xi_k\}_{k \ge 1}$ are i.i.d. then we have a compound Poisson process (time-changed random walk).
- first passage times for standard Brownian motion with drift.

Infinite divisibility and Lévy processes

There is a bijection between the class of Lévy processes and the class of infinitely divisible distributions.

Lévy processes \longleftrightarrow infinitely divisible distributions

For any
$$t \ge 0$$
 and any $n \in \mathbb{N}$,

$$X_t = \left(X_{\frac{tn}{n}} - X_{\frac{(t-1)n}{n}}\right) + \dots + \left(X_{\frac{1}{n}} - X_0\right).$$

• Furthermore, setting $\mu = \mathscr{L}(X_1)$, we have $\mathscr{L}(X_t) = \mu^{t*}$ (!). Therefore we can either talk of a particular Lévy process or of an infinitely divisible distribution, say μ .

Lévy-Khintchine formula

The main subject of this talk is the beautiful and fundamental,

Theorem (Lévy, Khintchine)

Let μ be an infinitely divisible distribution supported on \mathbb{R} . Then for any $\theta \in \mathbb{R}$ its characteristic function is of the form,

$$\widehat{\mu}(\theta) = \exp\left[ia\theta - \frac{1}{2}\sigma^2\theta^2 + \int_{\mathbb{R}} \left(e^{i\theta x} - 1 - i\theta x \mathbf{1}_{|x|<1}\right)\nu(dx)\right],$$

where $a, \sigma \in \mathbb{R}$ and ν is a measure satisfying,

$$u(\{0\}) = 0 \quad \text{and} \quad \int_{\mathbb{R}} (1 \wedge x^2) \nu(dx) < \infty.$$

We call (a, σ, ν) the characteristic triplet of μ , and ν is referred to as the Lévy measure (or sometimes the jump measure).

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Comments

- The quantity in the square brackets is called the characteristic exponent, ψ(θ).
- ► The cutoff function can be modified from 1_{|x|<1} but the drift term has to be appropriately modified.
- Notice that,

$$\left|\int_{|x|<1}(e^{i\theta x}-1-i\theta x)\nu(dx)\right|\leq Cu\int_{|x|<1}x^2\nu(dx)<\infty,$$

and also,

$$\left|\int_{|x|\geq 1} (e^{i\theta x}-1)\nu(dx)\right|\leq 2\int_{|x|\geq 1}\nu(dx)<\infty.$$

So the integrability conditions of the Lévy measure are quite natural.

Examples of characteristic exponents

▶ For $\mu \sim \mathcal{N}(\mathbf{a}, \sigma^2)$ (standard Brownian motion with drift) then,

$$\psi(\theta) = ia\theta - \frac{1}{2}\sigma^2\theta^2.$$

Clearly $\nu = 0$.

• If $\mu \sim \text{Poiss}(\lambda)$ for $\lambda > 0$ (Poisson process) then,

$$\psi(\theta) = \lambda(e^{i\theta} - 1).$$

Notice that $\nu = \lambda \delta_1$.

• When $\mu \sim \text{ComPoiss}(\lambda, F)$ we get

$$\psi(\theta) = \lambda \int_{\mathbb{R}} (e^{i\theta x} - 1)F(dx).$$

Observe that $a = -\lambda \int_{|x|<1} xF(dx)$ while $\nu(dx) = \lambda F(dx)$.

Lévy-Khintchine re-visited

One way to re-write the characeristic exponent is as follows,

$$\psi(\theta) = -\left[\frac{1}{2}\sigma^2\theta^2\right] + \left[\int_{\mathbb{R}} (e^{i\theta x} - 1)\nu(dx)\right] + \left[i\theta\left(a - \int_{|x| < 1} x\nu(dx)\right)\right].$$

Another way to re-write the exponent has a more probabilistic interpretation, and is related to the Lévy-Itō decomposition,

$$\psi(\theta) = \left[i\theta a - \frac{1}{2}\sigma^2\theta^2\right] + \left[\int_{|x|\geq 1} (e^{i\theta x} - 1)\nu(dx)\right] + \left[\int_{|x|<1} (e^{i\theta x} - 1 - i\theta x)\nu(dx)\right].$$

The proof of this decomposition gives deeper insight into the origins of the Lévy measure.

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Origin of the Lévy measure

Outline of the proof:

- First, for any Borel set, A, bounded away from zero one can define the counting process associated to the jump process of a Lévy process, counting jumps up to time t taking values in A (recall càdlàg).
- The resulting process is a Poisson process and gives rise to a Poisson random measure.
- ► The Lévy measure is the intensity measure of the Poisson random measure, when A = ℝ \ (-1, 1).
- ► Therefore the next step is to subtract, from any given Lévy process, the portion corresponding to jumps of size ≥ 1.
- Centering this process gives rise to an L²-martingale, which can be shown to split into the sum of a Brownian motion and a pure jump process.
- The square summability of the L² pure jump martingale together with the choice of A give the integrability conditions of ν.

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